

A note on pricing forwards in a regime-switching model via integro-PDE method

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ABSTRACT.

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1

1. MOTIVATION AND INTRODUCTION

2 In Geman and Roncoroni (2006) a Markov model for the spot price dynamics of electricity
3 is proposed. This model is a standard Ornstein-Uhlenbeck process, with a non-standard
4 jump term. The jumps are designed to model the spikes frequently observed for electricity
5 spot prices.

6 This threshold model is a Markov jump-diffusion model, however, not feasible for explicit
7 pricing of forward contracts due to its specification. A forward contract is an agreement
8 where the buyer purchases any specified commodity at an agreed time to an agreed price.
9 The agreed price, commonly known as the *forward price* is the price such that the current
10 value of the contract is zero. With a pricing measure it can be viewed as the best predicted
11 spot time at time of the transaction. Mathematically it is expressed as the conditional
12 expected spot price (possibly under a risk-neutral probability).

13 The problem we focus here is the derivation of the forward price dynamics. Unlike many
14 other models, the threshold model does not allow for explicit calculation of the forward
15 price, and numerical methods are called for. Since we want to find the dynamics, the Monte
16 Carlo method is very cumbersome, and we are going to analyse a PDE-based approach.
17 Since the threshold model involves jumps, we are led to integro-PDEs, and numerical
18 methods for such. We want to compare the resulting forward prices with those of similar
19 models.

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21 **2.1. Technical preliminaries.** In this section we recall some definitions and helpful facts
 22 we use in further sections. We mainly consider a Lévy-type process for modelling the
 23 power price. Let $(\Omega, \mathbb{P}, \mathcal{F}, \{\mathcal{F}_t\})$ be a complete filtered probability space. Let T defines the
 24 time horizon. We will use the following definitions given in Chapter 3 in Cont and Tankov
 25 (2004).

26 **Definition 1.** A cádlág stochastic process $(L_t)_{t \geq 0}$ on $(\Omega, \mathbb{P}, \mathcal{F}, \{\mathcal{F}_t\})$ with values in \mathbb{R} such
 27 that $L_0 = 0$ is called a Lévy process if it satisfies the following properties:

- 28 (1) *Independent increments:* for every increasing sequence of times t_0, \dots, t_n the ran-
 29 dom variables $X_{t_0}, X_{t_1} - X_{t_0}, \dots, X_{t_n} - X_{t_{n-1}}$ are independent;
- 30 (2) *Stationary increments:* the law of $X_{t+h} - X_t$ does not depend on t ;
- 31 (3) *Stochastic continuity:* $\forall \epsilon > 0, \lim_{h \rightarrow 0} \mathbb{P}(|X_{t+h} - X_t| \geq \epsilon) = 0$.

32 A Lévy process is associated with its Lévy measure ν

33 **Definition 2.** Let $(L_t)_{t \geq 0}$ be a Lévy process on \mathbb{R} . The measure ν on \mathbb{R} defined by:

$$(1) \quad \nu(A) = \mathbb{E} [\#\{t \in [0, 1] : \Delta L_t \neq 0, \Delta L_t \in A\}], \quad A \in \mathcal{B}(\mathbb{R})$$

34 is called the Lévy measure of L : $\nu(A)$ is the expected number, per unit time, of jumps whose
 35 size belongs to A .

36 Every Lévy process is characterised by its characteristic triplet (γ, b, ν) , where $\gamma \in \mathbb{R}$ is
 37 the drift term, $b \in \mathbb{R}_{\geq 0}$ is the diffusion coefficient and ν is the Lévy measure. Next useful
 38 theorem is a celebrated Lévy-Khincin decomposition.

39 **Theorem 1.** Let $(L_t)_{t \geq 0}$ be a Lévy process on \mathbb{R} with a characteristic triplet (γ, b, ν) . Then

$$(2) \quad \mathbb{E} [e^{izLt}] = e^{t\psi(z)}, \quad z \in \mathbb{R},$$

40 with

$$(3) \quad \psi(z) = i\gamma z - \frac{z^2 b}{2} + \int_{\mathbb{R}} (e^{izx} - 1 - izx \mathbf{1}_{|x| < 1}) \nu(dx).$$

41

42 *Proof.* The proof can be found in various sources, for example in Chapter 3 in Cont and
 43 Tankov (2004) and in Chapter 4 in Sato (1999). \square

44 According to this theorem a Lévy process can always be decomposed into several processes:
 45 a deterministic linear process (drift) with parameter γ , a Brownian motion with coefficient
 46 \sqrt{b} , a compound Poisson process with arrival rate $\lambda := \nu(\mathbb{R} \setminus (-1, 1))$ and jump size dis-
 47 tribution given by its cumulative distribution function $\mathbb{F}(dx) := \frac{\nu(dx)}{\nu(\mathbb{R} \setminus (-1, 1))} \mathbf{1}_{|x| \geq 1}$, and the
 48 last component: pure jump martingale process.

49 Now let us recall a few facts about the Lévy measure ν :

- 50 • the Lévy measure ν on \mathbb{R} satisfies $\nu(\{0\}) = 0$ and $\int_{\mathbb{R}} (1 \wedge (|x|)^2) \nu(dx) < \infty$.

- 51 • if ν is a finite measure, i.e. $\lambda = \nu(\mathbb{R}) = \int_{\mathbb{R}} \nu(dx) < \infty$, then $\mathbb{F}(dx) := \frac{\nu(dx)}{\lambda}$ is a
 52 probability measure. Then λ is interpreted as the expected number of jumps and
 53 $\mathbb{F}(dx)$ is the distribution of the jump size x . It is also said the the Lévy process L_t
 54 has finite activity. (Theorem 21.3 in Sato (1999)).
- 55 • if $b \neq 0$ or $\int_{|x| \leq 1} |x| \nu(dx) = \infty$, then almost all the paths of the Lévy process L_t
 56 have infinite variation. (Theorem 21.9 in Sato (1999)).

57 **Proposition 2.1.** *Let $(L_t)_{t \geq 0}$ be a Lévy process on \mathbb{R} with a characteristic triplet (γ, b, ν) .*
 58 *Then*

- 59 (1) $\mathbb{E}[|L_t|^p] < \infty$ if and only if $\int_{|x| \geq 1} |x|^p \nu(dx) < \infty$;
- 60 (2) $\mathbb{E}[e^{pL_t}] < \infty$ if and only if $\int_{|x| \geq 1} e^{px} \nu(dx) < \infty$.

61 *Proof.* Theorem 25.3 in Sato (1999). □

62 We now continue with a few more known facts and properties of a Lévy process and its
 63 connection to the martingale theory.

- a *semimartingale* is a stochastic process $(X_t)_{0 \leq t \leq T}$ that can be represented as

$$X = X_0 + M + A,$$

64 where X_0 is finite and \mathcal{F} -measurable, M is a local martingale with $M_0 = 0$ and A
 65 is a finite variation process with $A_0 = 0$;

- a semimartingale X is a *special semimartingale*, if the process A is predictable;
- every Lévy process is a semimartingale due to its Lévy-Khincin decomposition;
- every Lévy process with its finite first moment (i.e. if and only if $\int_{|x| \geq 1} |x| \nu(dx) < \infty$;) is also a special semimartingale;
- the following three assertions are equivalent:

- 71 (1) a Lévy process L_t is a special semimartingale;
- 72 (2) $\int_{\mathbb{R}} (|x| \wedge |x|^2) \nu(dx) < \infty$,
- 73 (3) $\int_{\mathbb{R}} (|x| \mathbf{1}_{|x| \geq 1}) \nu(dx) < \infty$,

74 this is a consequence of Lemma 2.8 in Kallsen and Shiryaev (2002).

Theorem 2. *Let $(L_t)_{t \geq 0}$ be a Lévy process on \mathbb{R} . Define*

$$L_t^* = \sup_{s \in [0, t]} |L_s|.$$

75 *Let $g(r)$ be a nonnegative continuous submultiplicative function on $[0, \infty)$, increasing to ∞*
 76 *as $r \rightarrow \infty$. Then the following four statements are equivalent:*

- 77 (1) $\mathbb{E}[g(L_t^*)] < \infty$ for some $t > 0$;
- 78 (2) $\mathbb{E}[g(L_t^*)] < \infty$ for every $t > 0$;
- 79 (3) $\mathbb{E}[g(|L_t^*|)] < \infty$ for some $t > 0$;

80 (4) $\mathbb{E}[g(|L_t^*|)] < \infty$ for every $t > 0$.

81 *Proof.* Theorem 25.18 in Sato (1999). □

82 **2.2. Power price modelling.** Let $S(t)$ be the spot price of electricity (power) defined
83 as

$$(4) \quad S(t) = \exp(\mu(t) + Y(t)),$$

84 where $\mu(t)$ is some deterministic seasonality function and $Y(t)$ is some stochastic process.
85 The classical approach to model the process $Y(t)$, as for example considered in Carlea and
86 Figueroa (2005), is stated as

$$(5) \quad dY(t) = -\alpha Y(t) dt + \sigma dW_t + dQ_t,$$

87 where W_t is a Brownian motion, Q_t is a compound Poisson process, α is the speed of
88 mean-reversion and σ is the volatility. Process W_t and Q_t are assumed to be mutually
89 independent.

90 Alternatively, Geman and Roncoroni (2006) proposed to model this stochastic component
91 differently

$$(6) \quad dX(t) = -\alpha X(t) dt + \sigma dW_t + h(X(t)) dQ_t,$$

92 where all the notations and assumptions are the same as for the process $Y(t)$, and $h(x)$ is
93 a state-dependent function which is -1 for large values of X (defined by some threshold
94 \mathcal{T}) and 1 otherwise. Despite the "regime-switching" term $h(x)$, this process holds the
95 Markov property in a single state variable, for a proof see Roncoroni (2002). The authors
96 claim that the process X_t is a special semimartingale. This model is referred here as the
97 threshold model.

98 The difference of the two models lies in the change of sign of the h -function. This function
99 ensures that the price may jump downwards in the case of high spot prices. Note that
100 when we use a minus in front of the jumps in the classical model, we want to use it in the
101 "high price" regime for the threshold model. By the "high price" regime we mean here
102 that the price is far above its mean level, which can happen when a spike or big jump
103 occurred.

104 We also notice here that the process $L_t := \sigma W_t + Q_t$ is a Lévy process in a contrast to
105 the process $L_t^h := \sigma W_t + h(x) Q_t$ which does not satisfy the properties given in Definition
106 1.

107 **2.3. Forward modelling.** We know that the forward price $F(t, T)$ at time t , for a contract
 108 with a delivery at time $T \geq t$, is

$$(7) \quad F(t, T) = \mathbb{E}^{\mathbb{Q}}[S(T) | \mathcal{F}_t],$$

109 which is a martingale under the equivalent martingale measure \mathbb{Q} . However, the power
 110 spot price does not need to be a martingale with respect to \mathbb{Q} since power is a non-
 111 storable commodity. So we call the measure \mathbb{Q} a pricing measure, as it is a probability
 112 that takes into account all the risks associated with the change in the price (spikes can,
 113 for example, happen due to sudden weather change or unexpected outage of equipment).
 114 The choice of \mathbb{Q} can be done via a canonical Girsanov (drift part) and Esscher (jump part)
 115 transformations. Alternatively, one could say that the process S_t is already under the
 116 measure \mathbb{Q} and one could argue that the market will charge an additional risk premium by
 117 changing/adjusting the mean level. This would mean the $\mathbb{Q} = \mathbb{P}$ with the latter being the
 118 real-world pricing measure. From now on we apply this assumption.

119 By the Markovian property we can write the forward price explicitly as a function of $X(t)$
 120 as

$$F(t, T, X(t)) = e^{\mu(T)} \mathbb{E}[e^{X(T)} | X(t)] = e^{\mu(T)} f(t, X(t)),$$

121 with $f(t, X(t)) := \mathbb{E}[e^{X(T)} | X(t) = x]$. Our aim here is to derive efficient routines to
 122 calculate the function $f(t, X(t))$ based on the associated integro-PDE in terms of the
 123 threshold model and to study the impact of the function $h(x)$ on the forward prices. We
 124 will also compare obtained forwards with ones from the standard classical model given in
 125 Equation (5).

126 **2.4. Forward price of the classical model.** This model allows for explicit forward price
 127 formula. More precisely, let the forward price be

$$(8) \quad G(t, T, Y(t)) = e^{\mu(T)} \mathbb{E}[e^{Y(T)} | Y(t)] = e^{\mu(T)} g(t, Y(t)),$$

128 with $g(t, Y(t)) \equiv \mathbb{E}[e^{Y(T)} | Y(t) = y]$. One can calculate the function $g(t, Y(t))$ analytically
 129 by appealing to the moment generating function of the compound Poisson process Q_t . We
 130 start with the dynamics of the logarithm of the price S_t

$$(9) \quad \begin{aligned} d \log S_t &= \mu'_t dt - \alpha (\log S_t - \mu(t)) dt + \sigma dW_t + dQ_t \\ &= \alpha (\hat{\mu}(t) - \log S_t) dt + \sigma dW_t + dQ_t, \end{aligned}$$

131 where $\hat{\mu}(t) := \frac{1}{\alpha} \mu'_t + \mu(t)$. Let us now apply Ito's lemma to $(e^{\alpha t} \log S_t)$ to obtain

$$(10) \quad d(e^{\alpha t} \log S_t) = \alpha e^{\alpha t} \hat{\mu}(t) dt + e^{\alpha t} \sigma dW_t + e^{\alpha t} dQ_t.$$

132 After integrating from t to T and replacing terms, we have

$$(11) \quad \log S_T = \mu(T) + Y(t)e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} dQ_s,$$

133 then the price S_T becomes

$$(12) \quad S_T = e^{\mu(T)+Y(T)} = e^{\mu(T)+Y(t)e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} dQ_s}.$$

134 Now coming back to Equation (8) we have that function $g(t, Y(t))$ becomes

$$(13) \quad \begin{aligned} g(t, Y(t)) &\equiv \mathbb{E}[\exp(Y(T)) \mid Y(t) = y] \\ &= \mathbb{E}[e^{Y(t)e^{-\alpha(T-t)}} e^{\sigma \int_t^T e^{-\alpha(T-s)} dW_s} e^{\int_t^T J e^{-\alpha(T-s)} dN_s} \mid Y(t) = y] \\ &= e^{ye^{-\alpha(T-t)}} \exp\left(\frac{\sigma^2}{4\alpha}(1 - e^{-2\alpha(T-t)})\right) \exp\left(\lambda \int_t^T (\mathbb{E}[e^{J e^{-\alpha(T-s)}}] - 1) ds\right), \end{aligned}$$

135 where for the last equality we use mutual independence of W_t and Q_t , the fact the
136 $\sigma \int_t^T e^{-\alpha(T-s)} dW_s$ is a normally distributed random variable and the Lévy-Khincin rep-
137 resentation for the compound Poisson process Q_t .

138 2.4.1. *Normal distribution for jumps.* Now following Cartea and Figueroa (2005), we as-
139 sume Normal distribution for the jump size J , i.e. $J \sim \mathcal{N}(m_1, m_2)$ with mean m_1 and
140 standard deviation m_2 . This allows us to compute

$$(14) \quad \mathbb{E}[e^{J e^{-\alpha(T-s)}}] = \exp\left(m_1 e^{-\alpha(T-s)} + \frac{m_2^2}{2} e^{-2\alpha(T-s)}\right).$$

141 Then the forward price $g(t, Y(t))$ when the jump size J follows Normal distribution is given
142 as

$$(15) \quad g(t, Y(t)) = e^{ye^{-\alpha(T-t)}} \exp\left(\frac{\sigma^2}{4\alpha}(1 - e^{-2\alpha(T-t)})\right) \exp\left(\lambda \int_t^T e^{m_1 e^{-\alpha(T-s)} + \frac{m_2^2}{2} e^{-2\alpha(T-s)}} ds - \lambda(T-t)\right).$$

143 2.4.2. *Laplace distribution for jumps.* Another option for the jump size distribution is a
144 Laplace distribution (see Section 5 for the details). Besides explanatory advantages such
145 as capturing the heavy tails spike nature, we have the following useful property: when
146 $J \sim \mathcal{Laplace}(m_1, m_2)$ with m_1 – the location parameter and $m_2 > 0$ – the scale parameter,
147 then $e^{cJ} \sim \mathcal{LogLaplace}(cm_1, cm_2)$ with some constant c . A very detailed investigation of
148 the LogLaplace distribution can be found in a book of Kozubowski and Podgorski (2003).
149 We can use their formula for the expected value of a random variable e^{cJ} and obtain

$$(16) \quad \mathbb{E}[e^{cJ}] = \frac{\delta}{1 - c^2 m_2^2},$$

150 where $\delta := e^{cm_1}$. Then this allows us to compute the expected value

$$(17) \quad \mathbb{E}[e^{Je^{-\alpha(T-s)}}] = \frac{\exp(e^{-\alpha(T-s)}m_1)}{1 - e^{-2\alpha(T-s)}m_2^2}.$$

151 So the the forward price $g(t, Y(t))$ when the jump size J follows Laplace distribution is
152 given as

$$(18) \quad g(t, Y(t)) = e^{ye^{-\alpha(T-t)}} \exp\left(\frac{\sigma^2}{4\alpha}(1 - e^{-2\alpha(T-t)})\right) \exp\left(\lambda \int_t^T \left(\frac{\exp(e^{-\alpha(T-s)}m_1)}{1 - e^{-2\alpha(T-s)}m_2^2}\right) ds - \lambda(T-t)\right).$$

153 **2.5. Forward price of the threshold model.** To get the forward price $f(t, X(t))$ with
154 the process $X(t)$ given in Equation (6) for the threshold model we start with calculating
155 the generator of the jump-diffusion process $X(t)$ for the threshold model. First, we know
156 that at time $t = T$ the forward price $F(T, T, X(T)) = S_T$ by the definition of the forward.
157 We can describe $f(t, X(t))$ by the backward Kolmogorov equation

$$(19) \quad f_t + \mathcal{L}f = 0, \text{ for } t < T, \text{ with } f(T, x) = e^x \text{ at } t = T,$$

158 where

$$(20) \quad \mathcal{L}f = -\alpha x f_x + \frac{\sigma^2}{2} f_{xx} + \lambda h(x) \mathbb{E}[f(x + J, t) - f(x, t)],$$

159 where the expected value in the last term is over the probability distribution of jumps. Let
160 us show how to get this backward integro-PDE. Knowing that the compensated Poisson
161 process $\tilde{N}_t = N_t - \lambda t$ is a martingale, we start with applying Ito's formula to $f(t, X(t))$
162 which gives

$$\begin{aligned} f(T, X(T)) - f(t, x) &= \int_t^T (\sigma f_x)(s, X(s)) dW_s + \int_t^T (f_s - \alpha x f_x + \frac{\sigma^2}{2} f_{xx})(s, X(s)) ds \\ &+ \int_t^T h(X(s)) [f(s, X(s) + J) - f(s, X(s))] dN_s \\ &= \int_t^T (\sigma f_x)(s, X(s)) dW_s + \int_t^T (f_s - \alpha x f_x + \frac{\sigma^2}{2} f_{xx})(s, X(s)) ds \\ &+ \int_t^T \lambda h(X(s)) [f(s, X(s) + J) - f(s, X(s))] ds \\ &+ \int_t^T h(X(s)) [f(s, X(s) + J) - f(s, X(s))] d\tilde{N}_s. \end{aligned}$$

163 Now let us take the expected values from both sides and divide by $T - t$ which yields

$$\begin{aligned}
& \frac{\mathbb{E}[f(T, X(T))] - f(t, x)}{T - t} \\
= & \frac{\int_t^T \left\{ (f_s - \alpha x f_x + \frac{\sigma^2}{2} f_{xx})(s, X(s)) + \lambda \mathbb{E} \left[h(X(s)) (f(s, X(s) + J) - f(s, X(s))) \right] \right\} ds}{T - t},
\end{aligned}$$

164 then take the limit with $T - t \rightarrow 0$ and obtain

$$(21) \quad 0 = f_t - \alpha x f_x + \frac{\sigma^2}{2} f_{xx} + \lambda h(x) \mathbb{E}[f(t, x + J) - f(t, x)],$$

165 since we need that $\mathbb{E}[f(T, X(T)) | X(t) = x] = f(t, X(t))$ to be fulfilled. This is exactly
166 Equation (19).

167 In the next section we will implement the numerical scheme that solves this partial differ-
168 ential equation with an integral term to investigate the property of the function $h(x)$ on
169 the forward price.

170 Before to continue with the numerical investigation, we make one more clarifying step. Let
171 us rewrite the term of the expected jump in Equation (20) as

$$\begin{aligned}
(22) \quad \mathbb{E}[f(t, x + J) - f(t, x)] &= \int_{-\infty}^{\infty} (f(t, x + y) - f(t, x)) f_Y(y) dy \\
&\cong \int_{-\infty}^{\infty} \frac{(f(t, x + y) - f(t, x))}{dx} y f_Y(y) dy \\
&= f_x \underbrace{\int_{-\infty}^{\infty} y f_Y(y) dy}_{\cong C}.
\end{aligned}$$

172 with $f_Y(y)$ is a probability density function of the random jump size J . We mean here
173 that the integral part in this equation behaves like a gradient term, approximately giving
174 a rise to a second order differential operator as the right-hand side of Equation (23). So
175 then Equation (21) can be re-written as

$$(23) \quad f_t \cong (-\alpha x + \lambda h(x)C) f_x + \frac{\sigma^2}{2} f_{xx},$$

176 which clearly indicates that the term in front of f_x is of discontinuous nature. It also points
177 out the curvature due to $h(x)$ function. This "discontinuous" curvature can clearly be seen
178 in Figure 1.

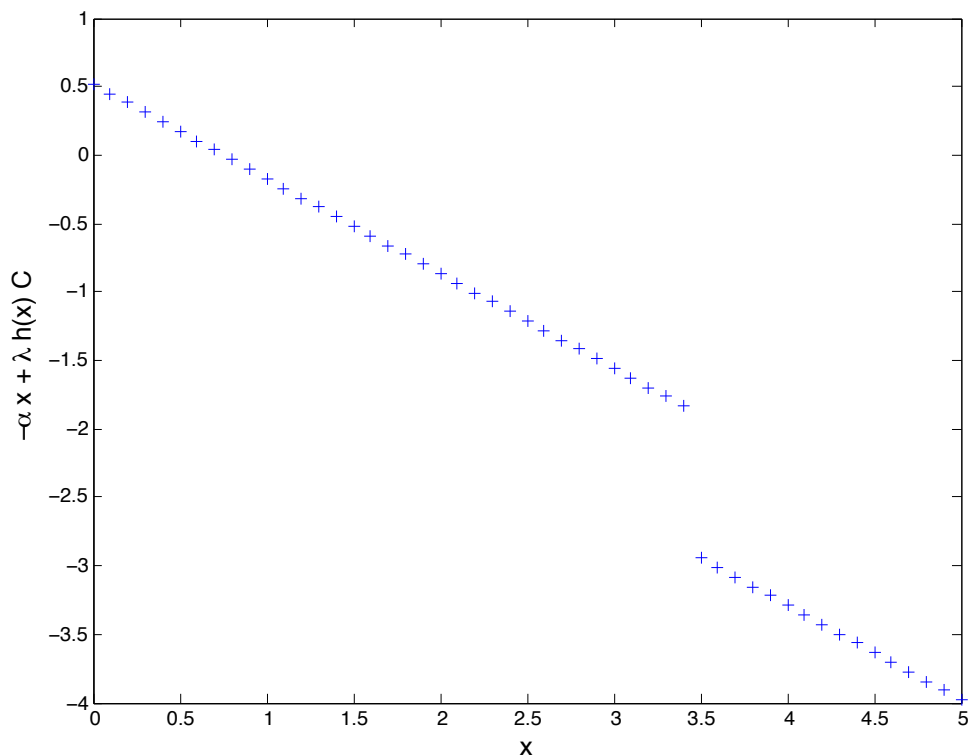


FIGURE 1. The curvature due to $h(x)$ function. Parameters: $C = 10$, $\alpha = 0.69$ (more than half a day), $\sigma = 2.59$, $\lambda = 13.5$ spikes per year, $\mathcal{T} = 3.5$.

179

3. NUMERICAL IMPLEMENTATION

180 This section contains three parts. Firstly, we discuss the method of finite difference which
 181 we apply to solve the integro-pde. When solving this equation we have to restrict our
 182 domain for the values x . It results in the truncation error which we discuss in the the
 183 second part of this section. Finally, we have to cut the integral term and investigate this
 184 truncation error as well.

185 **3.1. Method.** In this section we implement the finite difference method to solve equation
 186 (19). Excellent study of numerical methods in application to finance is given in the books
 187 of Cont and Tankov (2004) and Fusai and Roncoroni (2008). This scheme is an approxi-
 188 mation for this equation and is based on replacing derivatives by finite differences in the
 189 equation. The integral term responsible for the jumps is approximated by the Riemann
 190 sums. Since we have time and space variables we have to discretise the time and spacial
 191 domains.

192 Let us first re-write the Equation in (19) for the forward price $f(t, x)$ we solve via finite
 193 difference method

$$(24) \quad f_t = \alpha x f_x - \frac{\sigma^2}{2} f_{xx} + \lambda h(x) f - \lambda h(x) \int_{-\infty}^{\infty} f(t, x + y) f_Y(y) dy, \quad f(T, x) = e^x.$$

194 Then let us make the following replacement: $T - t = \tau$, which allows us to move backward
195 in time when solving the equation numerically. Then the integro-pde becomes

$$(25) \quad f_\tau = -\alpha x f_x + \frac{\sigma^2}{2} f_{xx} - \lambda h(x) f + \lambda h(x) \int_{-\infty}^{\infty} f(T - \tau, x + y) f_Y(y) dy, \quad f(\tau, x) = e^x.$$

196 Since there is enormous variety of sources on the finite difference method precisely applied
197 to solving financial mathematics problems, we will not focus on the details in general.
198 Instead, we provide the exact scheme we use here to solve Equation (25).

199 We start with a time domain for $\tau \in [0, T]$ and discretise it with $\Delta\tau = \frac{\tau}{N}$ with N being
200 the number of time steps. Then we continue with the space domain (x_{min}, x_{max}) and
201 $\Delta x = \frac{x_{max} - x_{min}}{M}$ with M being the number of space steps. We also need to introduce some
202 values K_{min} and K_{max} responsible for the interval for the jump size. The time and space
203 derivatives become

$$(26) \quad \frac{\partial f}{\partial \tau} \approx \frac{f^{n+1} - f^n}{\Delta\tau},$$

$$(27) \quad \frac{\partial f}{\partial x} \approx \frac{f_{i+1}^{n+1} - f_i^{n+1}}{\Delta x},$$

$$(28) \quad \frac{\partial^2 f}{\partial x^2} \approx \frac{f_{i+1}^{n+1} - 2f_i^{n+1} + f_{i-1}^{n+1}}{(\Delta x)^2},$$

204 with $n = 1, \dots, N + 1$ and $i = 1, \dots, M + 1$. Here we use a so-called explicit scheme.
205 Further we need to approximate the integral term. We do this via trapezoidal quadrature
206 rule with the same grid resolution Δx . As it is stated in Cont and Tankov (2004), due to
207 computational complexity when there is a jump term, it is more convenient to use implicit
208 scheme for the integral part

$$(29) \quad \begin{aligned} \int_{-\infty}^{\infty} f(T - \tau, x_i + y) f_Y(y) dy &\approx \lim_{\substack{K_{min} \rightarrow -\infty \\ K_{max} \rightarrow \infty}} \int_{K_{min}}^{K_{max}} f(T - \tau, x_i + y) f_Y(y) dy \\ &\approx \sum_{j=K_l}^{K_u} f^n(x_i + y_j) \int_{(j-1/2)\Delta x}^{(j+1/2)\Delta x} f_Y(y) dy \\ &= \sum_{j=K_l}^{K_u} f_{i+j}^n \int_{(j-1/2)\Delta x}^{(j+1/2)\Delta x} f_Y(y) dy, \end{aligned}$$

209 where $[K_{min}, K_{max}] \in [(K_l - 1/2)\Delta x, (K_l + 1/2)\Delta x]$.

210 So, the total scheme then becomes explicit-implicit, namely

$$(30) \quad \frac{f^{n+1} - f^n}{\Delta t} = Df^{n+1} + Jf^n,$$

211 where

$$(31) \quad (Df^{n+1})_i = \alpha x_i \frac{f_{i+1}^{n+1} - f_i^{n+1}}{\Delta x} - \frac{\sigma^2}{2} \frac{f_{i+1}^{n+1} - 2f_i^{n+1} + f_{i-1}^{n+1}}{(\Delta x)^2},$$

$$(32) \quad (Jf^n)_i = -\lambda h(x_i) \sum_{j=K_l}^{K_u} f_{i+j}^n \int_{(j-1/2)\Delta x}^{(j+1/2)\Delta x} f_Y(y) dy.$$

212 This results in

$$(33) \quad (I - \Delta\tau D)f^{n+1} = (I + \Delta\tau J)f^n,$$

213 where I is identity matrix, $(I - \Delta\tau D)$ and $(I + \Delta\tau J)$ are tridiagonal matrices. Since we
 214 know the final condition, i.e. $f(\tau, x) = e^x$ we move backward in time when solving this
 215 scheme. In other words, we know the f^n and search for the f^{n+1} .

216 As the reader can see, there are two things we should agree on when solving this equation:
 217 when $i = 1$ we need to know the term f_{i-1}^{n+1} which is out of the domain and when $i = M + 1$
 218 we need to know the term f_{i+1}^{n+1} which is also out of the domain. There are several possible
 219 solutions to this obstacle. One of them is to represent the second derivative differently
 220 and then not to go out of the domain of (x_{min}, x_{max}) . Another opportunity is to assume
 221 that our final condition at $\tau = 0$ is extended not only for the domain $[x_{min}, x_{max}]$ but for
 222 all the values x we need. However, in our case we can go for the third option and find
 223 these values explicitly since we have an exact solution for the forward price for the classical
 224 model considered above, at least for some distributions. We mean here the following: when
 225 $x > x_{max}$ it implies that the function $h(x) = -1$, which gives forward price $f(t, X(t))$ is
 226 almost equal to $g(t, Y(t))$: the difference is the minus sign in front of the jump component
 227 in Equation (5) (in the next section we show that if the jump size distribution is symmetric,
 228 then $f(t, X(t))$ is equal to $g(t, Y(t))$ when $x > x_{max}$ and $h(x) = -1$). When $x < x_{min}$ it
 229 implies that the function $h(x) = 1$, which gives forward price $f(t, X(t))$ is exactly equal to
 230 $g(t, Y(t))$.

231 There are two important points needed to be clarified when dealing with the finite difference
 232 scheme: consistence with the continuous equation and stability (i.e. the found solution
 233 does not blow up when $\Delta\tau$ and Δx goes to 0). The choice of those should be done
 234 carefully. The stability condition stated in Cont and Tankov (2004) and used here is
 235 $\Delta\tau \leq \inf\left\{\frac{1}{\lambda}, \frac{(\Delta x)^2}{\sigma^2}\right\}$.

236 **3.2. Domain truncation error.** When we solve numerically the partial differential equa-
 237 tion with an integral part, we have to define the domain for the x values. Defining this
 238 interval for $x \in (x_{min}, x_{max})$ means that we have to specify some boundary conditions at
 239 $x = x_{min}$ and $x = x_{max}$. On top of that we have specify the boundary values for our function

240 for the integral term. In our case we can specify these boundary values explicitly. When
 241 $x \leq x_{min}$, then the function $h(x) = 1$ and then the forward price $f(t, X(t)) = g(t, Y(t))$.
 242 When $x \geq x_{max}$, the function $h(x) = -1$, then the forward price $f(t, X(t)) = g^-(t, Y(t))$
 243 with minus in front of the jump component in Equation (5), namely

$$(34) \quad dY^-(t) = -\alpha Y(t) dt + \sigma dW_t - dQ_t.$$

244 Let us now compute the solution associated with this process $Y^-(t)$ assuming that the
 245 starting value y is the same as for the function $g(t, y)$

$$(35) \quad g^-(t, y) = e^{ye^{-\alpha(T-t)}} \exp\left(\frac{\sigma^2}{4\alpha}(1 - e^{-2\alpha(T-t)})\right) \exp\left(\lambda \int_t^T (\mathbb{E}[e^{-Je^{-\alpha(T-s)}}] - 1) ds\right).$$

246 We can see that the only difference between $g(t, y)$ and $g^-(t, y)$ lies in the expected values
 247 of the exponent of the jump component, let us compare these terms keeping in mind that
 248 $e^{-\alpha(T-s)} \in (0, 1]$ for $\alpha > 0$ and $T - s \geq 0$

$$\begin{aligned} & \mathbb{E}[e^{Je^{-\alpha(T-s)}} - e^{-Je^{-\alpha(T-s)}}] \\ = & \int_{-\infty}^{+\infty} \left(e^{xe^{-\alpha(T-s)}} - e^{-xe^{-\alpha(T-s)}} \right) f_X(x) dx \\ = & \int_{-\infty}^0 \left(e^{xe^{-\alpha(T-s)}} - e^{-xe^{-\alpha(T-s)}} \right) f_X(x) dx + \int_0^{+\infty} \left(e^{xe^{-\alpha(T-s)}} - e^{-xe^{-\alpha(T-s)}} \right) f_X(x) dx \\ = & \int_0^{+\infty} \left(e^{-xe^{-\alpha(T-s)}} - e^{xe^{-\alpha(T-s)}} \right) f_X(-x) dx + \int_0^{+\infty} \left(e^{xe^{-\alpha(T-s)}} - e^{-xe^{-\alpha(T-s)}} \right) f_X(x) dx \\ = & \begin{cases} 0, & \text{if distribution is symmetric} \\ \int_0^{+\infty} \left(e^{xe^{-\alpha(T-s)}} - e^{-xe^{-\alpha(T-s)}} \right) (f_X(x) - f_X(-x)) dx, & \text{otherwise} \end{cases} \\ = & \begin{cases} 0, & \text{if } f_X(x) = f_X(-x), \text{ distribution is symmetric;} \\ C_{>0}, & \text{if } f_X(x) > f_X(-x), \text{ upward jumps are more likely;} \\ C_{<0}, & \text{if } f_X(x) < f_X(-x), \text{ downward jumps are more likely.} \end{cases} \end{aligned}$$

(36)

249 So we observe that

$$\begin{cases} g(t, y) = g^-(t, y), & \text{distribution is symmetric;} \\ g(t, y) > g^-(t, y), & \text{upward jumps are more likely;} \\ g(t, y) < g^-(t, y), & \text{downward jumps are more likely;} \end{cases}$$

250 Let us then denote the solution to our boundary problem given in Equation (19) as

$$(37) \quad f_B(t, x) = \begin{cases} \hat{f}_B(t, x), & \text{if } x \in (x_{min}, x_{max}) \\ g(t, y), & \text{if } x \leq x_{min} \\ g^-(t, y), & \text{if } x \geq x_{max} \end{cases}$$

251 where $\hat{f}_B(t, x)$ solves Equation (24) on a bounded domain (x_{min}, x_{max}) , solution $g(t, y)$ is
 252 given in Equation (18) and solution $g^-(t, y)$ is given in Equation (35). Now we are ready
 253 to calculate the domain truncation error with the following proposition.

254 **Proposition 3.1.** *Assume that*

255 • for $\nu(dx)$ being a Lévy measure and for $\epsilon > 0$, $\int_{|x|>1} (e^{\epsilon|x|} - 1) \nu(dx) < \infty$;

256 • $x_{min} = -x_{max}$.

257 Let $f(t, x)$ be the solution of our problem in Equation (19) on the unbounded region and
 258 $f_B(t, x)$ be the solution defined in Equation (37). Then

$$|f(t, x) - f_B(t, x)| \leq 2C e^{-\epsilon(x_{max}-|\bar{x}|)} \begin{cases} g(t, y), & \text{if } f_X(x) \geq f_X(-x), \\ g^-(t, y), & \text{if } f_X(x) < f_X(-x), \end{cases}$$

259 where C is some constant and does not depend on x_{max} and $f_X(\cdot)$ is a probability density
 260 for a jump size distribution.

Proof. For $t < T$ denote by

$$Z(T) := \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} h(X(s)) dQ_s,$$

261 then $X(T) = xe^{-\alpha(T-t)} + Z(T)$. Denote also by $M_T^x := \sup_{t \leq s \leq T} |X(s)|$. Then for $\tau := \inf\{s \geq$
 262 $t \mid M_s^x \geq x_{max}\}$ being the first time when the process $X(t)$ leaves the interval (x_{min}, x_{max})
 263 we have

$$\begin{aligned} & |f(t, x) - f_B(t, x)| \\ &= \left| \mathbb{E} [e^{X(T)}] - \mathbb{E} [e^{X(T)} \mathbf{1}_{M_T^x < x_{max}}] - \mathbb{E} [e^{X(\tau)} \mathbf{1}_{M_T^x \geq x_{max}}] \right| \\ &= \left| \mathbb{E} [e^{X(T)} \mathbf{1}_{M_T^x \geq x_{max}}] - \mathbb{E} [e^{X(\tau)} \mathbf{1}_{M_T^x \geq x_{max}}] \right| \\ &\leq 2 \mathbb{E} [|e^{X(T)} \mathbf{1}_{M_T^x \geq x_{max}}|] \\ &= 2 \left(g^-(t, y) \mathbb{P} \left(\sup_{t \leq s \leq T} X(s) \geq x_{max} \right) + g(t, y) \mathbb{P} \left(\sup_{t \leq s \leq T} (-X(s)) \geq x_{max} \right) \right) \\ &= 2 \left(g^-(t, y) \mathbb{P} \left(\sup_{t \leq s \leq T} X(s) \geq x_{max} \right) + g(t, y) \mathbb{P} \left(- \inf_{t \leq s \leq T} X(s) \geq x_{max} \right) \right) \\ &= 2 \left(g^-(t, y) \mathbb{P} \left(\sup_{t \leq s \leq T} X(s) \geq x_{max} \right) + g(t, y) \mathbb{P} \left(\inf_{t \leq s \leq T} X(s) \leq x_{min} \right) \right) \end{aligned}$$

$$\begin{aligned}
&\leq \begin{cases} 2g(t, y) \left(\mathbb{P} \left(\sup_{t \leq s \leq T} X(s) \geq x_{max} \right) + \mathbb{P} \left(\inf_{t \leq s \leq T} X(s) \leq x_{min} \right) \right), & \text{if } f_X(x) \geq f_X(-x), \\ 2g^-(t, y) \left(\mathbb{P} \left(\sup_{t \leq s \leq T} X(s) \geq x_{max} \right) + \mathbb{P} \left(\inf_{t \leq s \leq T} X(s) \leq x_{min} \right) \right), & \text{if } f_X(x) < f_X(-x), \end{cases} \\
&= \begin{cases} 2g(t, y) \mathbb{P}(M_T^x \geq x_{max}), & \text{if } f_X(x) \geq f_X(-x), \\ 2g^-(t, y) \mathbb{P}(M_T^x \geq x_{max}), & \text{if } f_X(x) < f_X(-x). \end{cases} \\
&\quad (38)
\end{aligned}$$

264 Since the function $h(x)$ is bounded by -1 and 1 , we observe that

$$\begin{aligned}
Z(T) &= \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} h(X(s)) dQ_s \\
&\leq \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s \\
(39) \quad &=: \tilde{Z}^T,
\end{aligned}$$

265 where \tilde{Q}_t is a compound Poisson process and \tilde{Z}^T is a Lévy process with a measure $\nu(dx)$.

$$(40) \quad M_T^x = \sup_{t \leq s \leq T} |xe^{-\alpha(s-t)} + Z(s)| \leq \sup_{t \leq s \leq T} \left| \underbrace{xe^{-\alpha(s-t)}}_{=: \tilde{x}} + \tilde{Z}(s) \right| =: M_T^{\tilde{x}}.$$

266 We note that

$$(41) \quad \mathbb{P}(M_T^x \geq x_{max}) \leq \mathbb{P}(M_T^{\tilde{x}} \geq x_{max}).$$

267 Now with the help of Theorem 2 we find that

$$(42) \quad C := \mathbb{E}e^{\epsilon M_T^{\tilde{0}}} < \infty.$$

268 Applying further Chebyshev's inequality gives

$$(43) \quad \mathbb{P}(M_T^{\tilde{0}} \geq x_{max}) \leq Ce^{-\epsilon x_{max}}.$$

269 Now we show that

$$\begin{aligned}
\mathbb{P}(M_T^{\tilde{x}} \geq x_{max}) &= \mathbb{P} \left(\sup_{t \leq s \leq T} |\tilde{x} + \tilde{Z}(s)| \geq x_{max} \right) \\
&\leq \mathbb{P} \left(\sup_{t \leq s \leq T} |\tilde{Z}(s)| + |\tilde{x}| \geq x_{max} \right)
\end{aligned}$$

$$\begin{aligned}
&= \mathbb{P} \left(\sup_{t \leq s \leq T} |\tilde{Z}(s)| \geq x_{max} - |\tilde{x}| \right) \\
&= \mathbb{P} \left(M_T^{\tilde{0}} \geq x_{max} - |\tilde{x}| \right) \\
(44) \quad &\leq C e^{-\epsilon(x_{max} - |\tilde{x}|)}.
\end{aligned}$$

270 □

271 **3.3. Jump size domain truncation error.** Here we discuss the integral term truncation
272 error which we obtain when we cut the interval for the jump size as it was done in Equation
273 (29). Process Q_t responsible for the jump component in Equation (6) is a compound Poisson
274 process with a jump measure $\nu(dx)$ that measures the expected number of jumps per unit
275 time whose size belong to a set $A \in \mathcal{B}(\mathbb{R})$.

276 Now let us introduce a new compound Poisson process Q_t^K with a new jump measure
277 $\nu^K := \nu(dx) 1_{x \in [K_{min}, K_{max}]}$. Then we show the relation between the process $X(T)$ (with
278 the jump component formed by the process Q_t) for some $t \leq T$ assuming that $X(t) = x$
279 at time t and having that function $h(x)$ is bounded by -1 and 1

$$\begin{aligned}
X(T) &= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} h(X(s)) dQ_s \\
&= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} h(X(s)) \Delta Q(s) \\
&= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} h(X(s)) \sum_{i=N(t)+1}^{N(s)} J_i \\
&= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} h(X(s)) J \\
&\leq x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} |J| \\
&= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} \sum_{i=N(t)+1}^{N(s)} |J| \\
&= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \sum_{t < s \leq T} e^{-\alpha(T-s)} \Delta \tilde{Q}(s) \\
(45) \quad &= x e^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s,
\end{aligned}$$

280 where in the third equality we used a definition of a compound Poisson process given in Sec-
281 tion 11.3.1 in Shreve (2004). Analogously, the process $X^K(T)$ (with the jump component
282 formed by the process Q_t^K) with $X^K(t) = x$ at time t is

$$\begin{aligned}
X^K(T) &= xe^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} h(X^K(s)) dQ_s^K \\
(46) \quad &\leq xe^{-\alpha(T-t)} + \sigma \int_t^T e^{-\alpha(T-s)} dW_s + \int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K,
\end{aligned}$$

283 where J^K is a random variable responsible for the jump size and that falls into the interval
284 $I = [K_{min}, K_{max}]$.

285 Let us now introduce a solution $f_I(t, x)$ which solves the problem given in Equation (19)
286 but with the truncated jump size domain $[K_{min}, K_{max}]$ and a solution $f(t, x)$ with an
287 unbounded jump size domain $(-\infty, +\infty)$. We give an error estimate associated to this
288 truncation. This error estimate derivation is similar to the proof of Proposition 4.2 in Cont
289 and Voltchkova (2005). The main difference here is the presence of function $h(x)$ due to
290 which our process $X(t)$ is not a Lévy process.

291 **Proposition 3.2.** *Assume that a jump measure $\nu(dx)$ for a compound Poisson process*
292 *satisfies the following*

- 293 • $\int_{\mathbb{R}} \nu(dx) < \infty$;
- 294 • $K_{max} > 1$ and $K_{min} = -K_{max}$;
- 295 • for $\xi_1, \xi_2 > 0$, $\int_{-\infty}^{-1} |x| e^{\xi_1|x|} \nu(dx) < \infty$ and $\int_1^{\infty} |x| e^{\xi_2|x|} \nu(dx) < \infty$;
- 296 • for $\epsilon > 0$, $\int_{\mathbb{R}} (e^{\epsilon|x|} - 1) \nu(dx) < \infty$ (for the details see Cont and Tankov (2004),
297 Proposition 3.8.).

298 Then

$$\begin{aligned}
&|f(t, x) - f_I(t, x)| \\
&\leq e^{C_1^{t,T} + \frac{1}{2}C_1^{t,T}} e^{(T-t)(C_3 e^{-\beta_1|K_{min}|} + C_4 e^{-\beta_2|K_{max}|})} \left(e^{(T-t)(C_5 e^{-\beta_1|K_{min}|} + C_6 e^{-\beta_2|K_{max}|})} - 1 \right) \\
&+ 2(1 - e^{-\alpha(T-t)}) (C_7 e^{-\beta_1|K_{min}|} + C_8 e^{-\beta_2|K_{max}|}), \\
(47)
\end{aligned}$$

299 where $C_1^{t,T} = xe^{-\alpha(T-t)}$, $C_2^{t,T} = \frac{\sigma^2}{2\alpha}(1 - e^{-2\alpha(T-t)})$, some constants $\beta_1, \beta_2 > 0$ and come
300 constants C_3, C_4, C_5, C_6, C_7 and C_8 .

301 *Proof.* We recall a few useful facts and give some notations:

- 302 • $|e^x - 1| = (e^x - 1) + 2(1 - e^x)^+ \leq (e^x - 1) + 2|x|$;
- 303 • by $U_T := xe^{-\alpha(T-t)} + \int_t^T \sigma^{-\alpha(T-s)} dW_s$ we denote a process that follows Gaussian
304 distribution with mean $C_1^{t,T}$ and variance $C_2^{t,T}$;
- 305 • by \tilde{Q}_t we denote a compound Poisson process with a measure $\tilde{\nu}(dx) := \nu(dx) 1_{|x|}$;
- 306 • by \tilde{Q}_t^K we denote a compound Poisson process with a measure $\tilde{\nu}^K(dx) := \nu(dx) 1_{|x| \in [K_{min}, K_{max}]}$;

307 • by $D_t^K := \tilde{Q}_t - \tilde{Q}_t^K$ we denote a compound Poisson process with a measure $\hat{\nu}^K(dx) =$
 308 $\tilde{\nu}(dx) - \tilde{\nu}^K(dx) = \nu(dx) 1_{|x| \notin [K_{min}, K_{max}]}$;

309 Then we have

$$\begin{aligned}
 & |f(t, x) - f_I(t, x)| \\
 = & \left| \mathbb{E} \left[e^{U_T + \int_t^T e^{-\alpha(T-s)} h(X(s)) dQ_s} - e^{U_T + \int_t^T e^{-\alpha(T-s)} h(X^K(s)) dQ_s^K} \right] \right| \\
 \leq & e^{C_1^{t,T} + \frac{1}{2} C_2^{t,T}} \left| \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s} - e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K} \right] \right| \\
 = & e^{C_1^{t,T} + \frac{1}{2} C_2^{t,T}} \left| \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K} \left(e^{\int_t^T e^{-\alpha(T-s)} dD_s^K} - 1 \right) \right] \right| \\
 \leq & e^{C_1^{t,T} + \frac{1}{2} C_2^{t,T}} \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K} \right] \mathbb{E} \left[\left| e^{\int_t^T e^{-\alpha(T-s)} dD_s^K} - 1 \right| \right] \\
 \leq & e^{C_1^{t,T} + \frac{1}{2} C_2^{t,T}} \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K} \right] \left(\mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} dD_s^K} - 1 \right] + 2\mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dD_s^K \right| \right] \right).
 \end{aligned} \tag{48}$$

310 Now let us show that every term is bounded by some constant. Since the first term is
 311 easily computable, we start with the second term

$$\begin{aligned}
 & \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} d\tilde{Q}_s^K} \right] \\
 = & e^{\int_t^T \int_{\mathbb{R}} (e^{e^{-\alpha(T-s)}|x|} - 1) \tilde{\nu}^K(dx) ds} \\
 = & e^{\int_t^T \int_{K_{min}}^0 (e^{e^{-\alpha(T-s)}(-x)} - 1) \nu(dx) ds} + e^{\int_t^T \int_0^{K_{max}} (e^{e^{-\alpha(T-s)}x} - 1) \nu(dx) ds} \\
 \leq & e^{\int_t^T \int_{K_{min}}^0 (e^{-x} - 1) \nu(dx) ds} + e^{\int_t^T \int_0^{K_{max}} (e^x - 1) \nu(dx) ds} \\
 = & e^{\int_t^T \int_{\mathbb{R}} (e^{|x|} - 1) \tilde{\nu}^K(dx) ds} \\
 = & e^{(T-t)} \left(e^{-\beta_1 |K_{min}|} \int_{K_{min}}^0 (e^{|x| + \beta_1 |K_{min}| - e^{\beta_1 |K_{min}|}}) \nu(dx) + e^{-\beta_2 |K_{max}|} \int_0^{K_{max}} (e^{|x| + \beta_2 |K_{max}| - e^{\beta_2 |K_{max}|}}) \nu(dx) \right) \\
 \leq & e^{(T-t)} \left(e^{-\beta_1 |K_{min}|} \int_{K_{min}}^0 (e^{|x| + \beta_1 |K_{min}|} - 1) \nu(dx) + e^{-\beta_2 |K_{max}|} \int_0^{K_{max}} (e^{|x| + \beta_2 |K_{max}|} - 1) \nu(dx) \right) \\
 \leq & e^{(T-t)} (C_3 e^{-\beta_1 |K_{min}|} + C_4 e^{-\beta_2 |K_{max}|}).
 \end{aligned} \tag{49}$$

312 Analogously, we consider the first term in the sum in the brackets of Equation (48)

$$\begin{aligned}
 & \mathbb{E} \left[e^{\int_t^T e^{-\alpha(T-s)} dD_s^K} - 1 \right] \\
 = & e^{\int_t^T \int_{\mathbb{R}} (e^{e^{-\alpha(T-s)}|x|} - 1) \hat{\nu}^K(dx) ds} - 1 \\
 = & e^{\int_t^T \int_{-\infty}^{K_{min}} (e^{e^{-\alpha(T-s)}(-x)} - 1) \nu(dx) ds} + e^{\int_t^T \int_{K_{max}}^{+\infty} (e^{e^{-\alpha(T-s)}x} - 1) \nu(dx) ds} - 1 \\
 \leq & e^{\int_t^T \int_{-\infty}^{K_{min}} (e^{-x} - 1) \nu(dx) ds} + e^{\int_t^T \int_{K_{max}}^{+\infty} (e^x - 1) \nu(dx) ds} - 1 \\
 = & e^{\int_t^T \int_{\mathbb{R}} (e^{|x|} - 1) \hat{\nu}^K(dx) ds} - 1
 \end{aligned}$$

$$\begin{aligned}
&\leq e^{(T-t)} \left(e^{-\beta_1 |K_{min}|} \int_{-\infty}^{K_{min}} (e^{|x|+\beta_1 |K_{min}|-1}) \nu(dx) + e^{-\beta_2 |K_{max}|} \int_{K_{max}}^{+\infty} (e^{|x|+\beta_2 |K_{max}|-1}) \nu(dx) \right) - 1 \\
&\leq e^{(T-t)} (C_5 e^{-\beta_1 |K_{min}|} + C_6 e^{-\beta_2 |K_{max}|}) - 1.
\end{aligned}
\tag{50}$$

313 Finally we compute the boundary for $\mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dD_s^K \right| \right]$. This process D_t^K can be rep-
314 resented as a sum of two compound Poisson processes: P_t^K with a measure $\nu(dx) 1_{x > K_{max}}$
315 and N_t^K with a measure $\nu(dx) 1_{x < K_{min}}$. More precisely, the process $P_t^K \geq 0$ has only pos-
316 itive jumps not smaller than $K_{max} > 1$ and the process $N_t^K \leq 0$ has only negative jumps
317 not greater than $K_{min} < -1$. Then we have

$$\begin{aligned}
&\mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dD_s^K \right| \right] \\
&= \mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dP_s^K + \int_t^T e^{-\alpha(T-s)} dN_s^K \right| \right] \\
&\leq \mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dP_s^K \right| \right] + \mathbb{E} \left[\left| \int_t^T e^{-\alpha(T-s)} dN_s^K \right| \right] \\
&= \mathbb{E} \left[\int_t^T e^{-\alpha(T-s)} dP_s^K \right] - \mathbb{E} \left[\int_t^T e^{-\alpha(T-s)} dN_s^K \right] \\
&\leq \int_t^T e^{-\alpha(T-s)} \int_{-\infty}^{K_{min}} |x| \nu(dx) ds + \int_t^T e^{-\alpha(T-s)} \int_{K_{max}}^{+\infty} |x| \nu(dx) ds \\
&\leq (1 - e^{-\alpha(T-t)}) \left(e^{-\beta_1 |K_{min}|} \int_{-\infty}^{K_{min}} |x| e^{\beta_1 |K_{min}|} \nu(dx) + e^{-\beta_2 |K_{max}|} \int_{K_{max}}^{+\infty} |x| e^{\beta_2 |K_{max}|} \nu(dx) \right) \\
&\leq (1 - e^{-\alpha(T-t)}) (C_7 e^{-\beta_1 |K_{min}|} + C_8 e^{-\beta_2 |K_{max}|}).
\end{aligned}
\tag{51}$$

318 □

319 4. RESULTS AND DISCUSSION

320 **4.1. Threshold and Cartea model forwards.** This section demonstrates the results
321 of two models. We plot the forward prices for the threshold and for the classical model.
322 For both models for the sake of simplicity we assume zero seasonal component $\mu(t)$. We
323 take some calibrated parameters obtained from the German spot power market (see pa-
324 pers Benth, Kiesel, and Nazarova (2012) and Bannoer, Kiesel, Nazarova, and Scherer
325 (2013)). Table 1 contains all the estimated parameters we use.

326 Here we provide the forward prices $f(t, X(t))$ as a solution to Equation (25) together with
327 the forward prices given in Equation (15).

328 We start with assumption that jump size random variables (log scale) follow Gaussian
329 distribution. Figure 4 shows prices $f(t, X(t))$ and $g(t, Y(t))$. When we are at the maturity,
330 i.e. $T - t = 0$, we observe our boundary condition. We also can see that the prices decrease

TABLE 1. An overview over the estimated parameter values for the forward price.

Parameter	Interpretation	Estimated value	Measure unit
α	mean-reversion force	0.6923	approx. 1 day
σ	volatility	2.59	
λ	jump intensity	13.5	spikes per year
\mathcal{T}	jump threshold	3	log scale
K_{max}	jump truncation	8	log scale

331 when time to maturity increases. We observe that both models produce similar results, in
 332 general the price level is almost the same. However, the threshold model produces slightly
 333 lower prices compared to the classical approach. This is perfectly in line with the behavior
 334 of function $h(x)$.

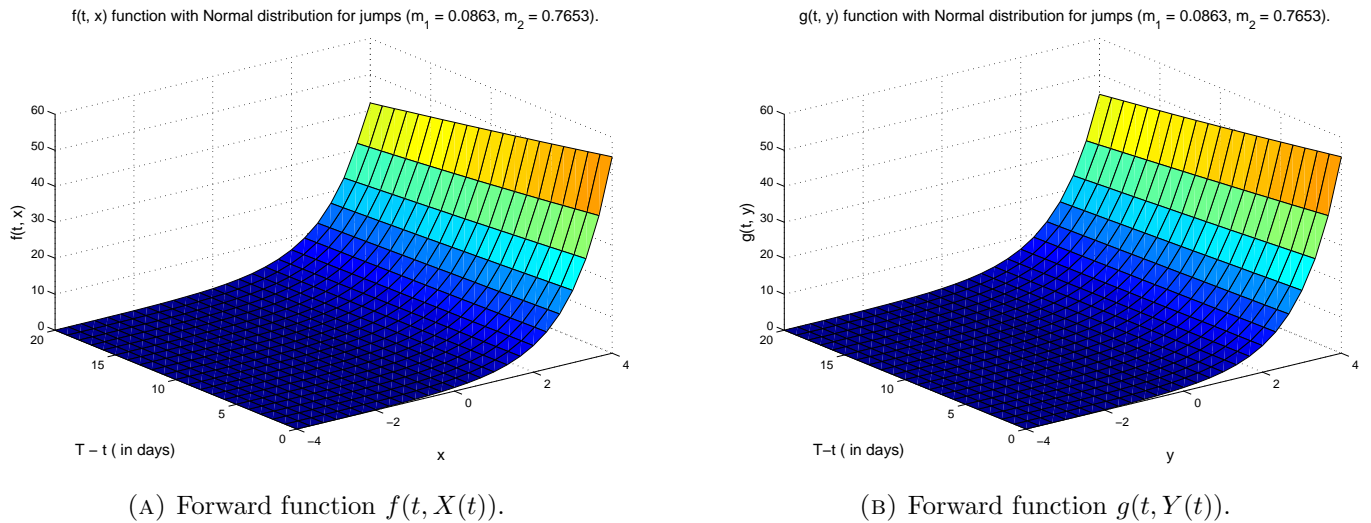
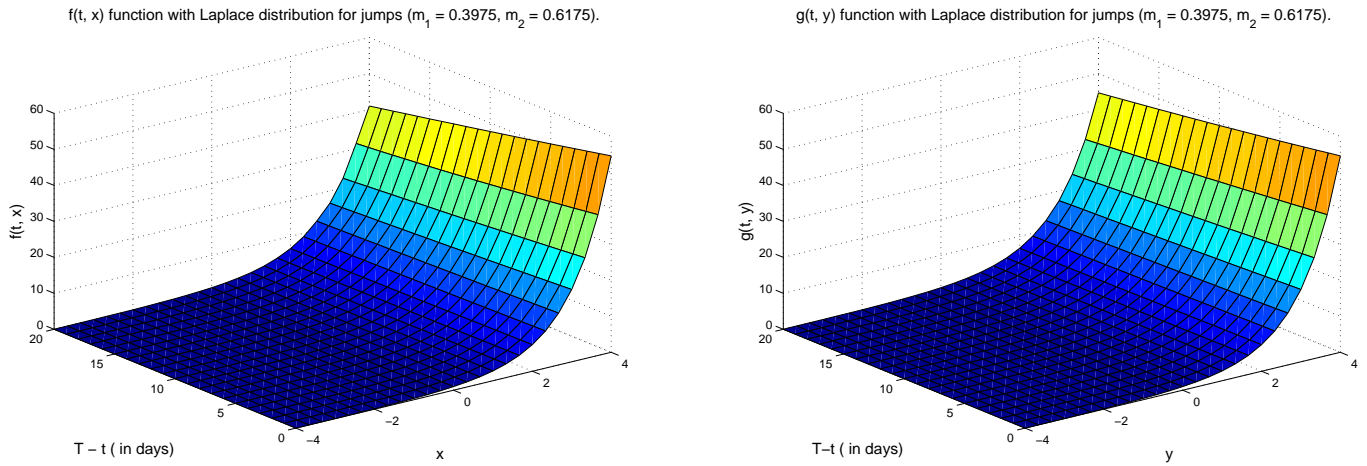


FIGURE 2. Normal distribution assumed for jumps with mean $m_1 = 0.0863$ and standard deviation $m_2 = 0.7653$. Time to maturity $T - t = 20$ days.

335 We continue with assumption that jump size random variables (log scale) follow Laplace
 336 distribution. Figure 5 shows prices $f(t, X(t))$ and $g(t, Y(t))$. When we are at the maturity,
 337 i.e. $T - t = 0$, we again observe our boundary condition fulfilled. We also can see that the
 338 prices decrease when time to maturity increases. Again both models demonstrate similar
 339 results. Moreover, different jump size distribution does not provide a significant difference
 340 to the price level. Again, the threshold model produces slightly lower prices compared to
 341 the classical approach due to the function $h(x)$.

342 BIG AVERAGE JUMP SIZE

343 4.2. **Discussion.** The function h does not show a remarkable impact on the forward prices.
 344 The prices are slightly smaller compared to those from the classical model.



(A) Forward function $f(t, X(t))$ in the log scale.

(B) Forward function $g(t, Y(t))$ in the log scale.

FIGURE 3. Laplace distribution assumed for jumps with location $m_1 = 0.3975$ and scale $m_2 = 0.6175$. Time to maturity $T - t = 20$ days.

TABLE 2. Comparative final forward values for the threshold ($\mathcal{T} = 3$) and jump-diffusion models for various starting values x and various average jump size parameter m_1 values. $\tau = 20$ days to maturity.

		x = y = 4		x = y = 2	
		$f(\tau, x)$	$g(\tau, y)$	$f(\tau, x)$	$g(\tau, y)$
Normal	$m_1 = 0.0863$	42.0727	44.4889	6.5248	6.5248
	$m_1 = 0.27$	42.0345	44.5389	6.5254	6.5248
Laplace	$m_1 = 0.3975$	40.9618	44.6442	6.6176	6.5248
	$m_1 = 1.2$	37.9248	45.1348	6.8678	6.5248

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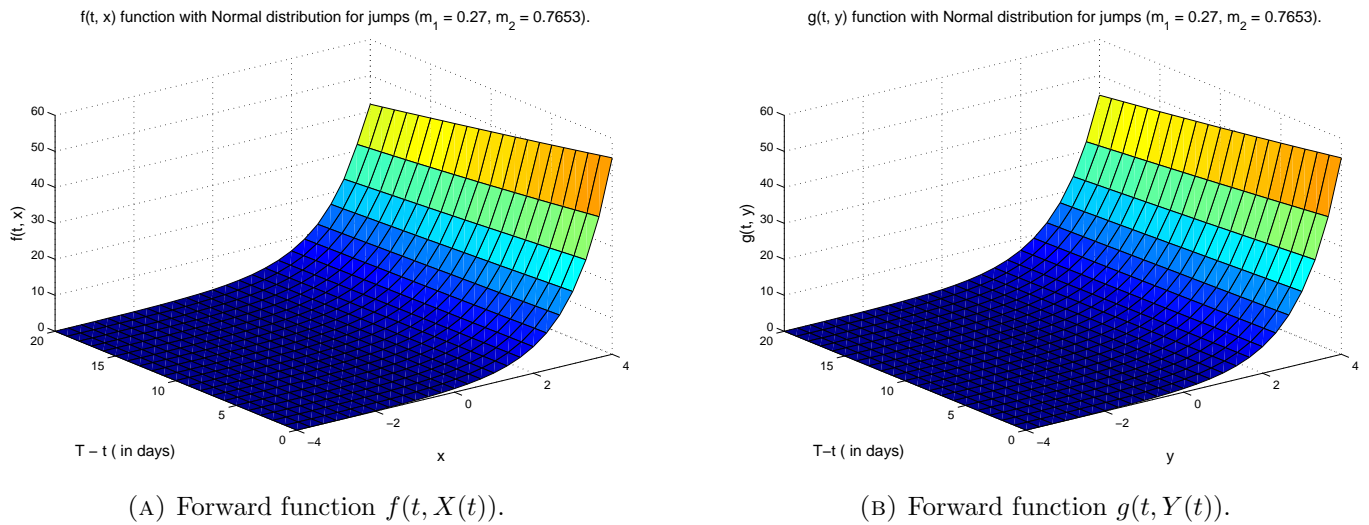


FIGURE 4. Normal distribution assumed for jumps with mean $m_1 = 0.27$ and standard deviation $m_2 = 0.7653$. Time to maturity $T - t = 20$ days.

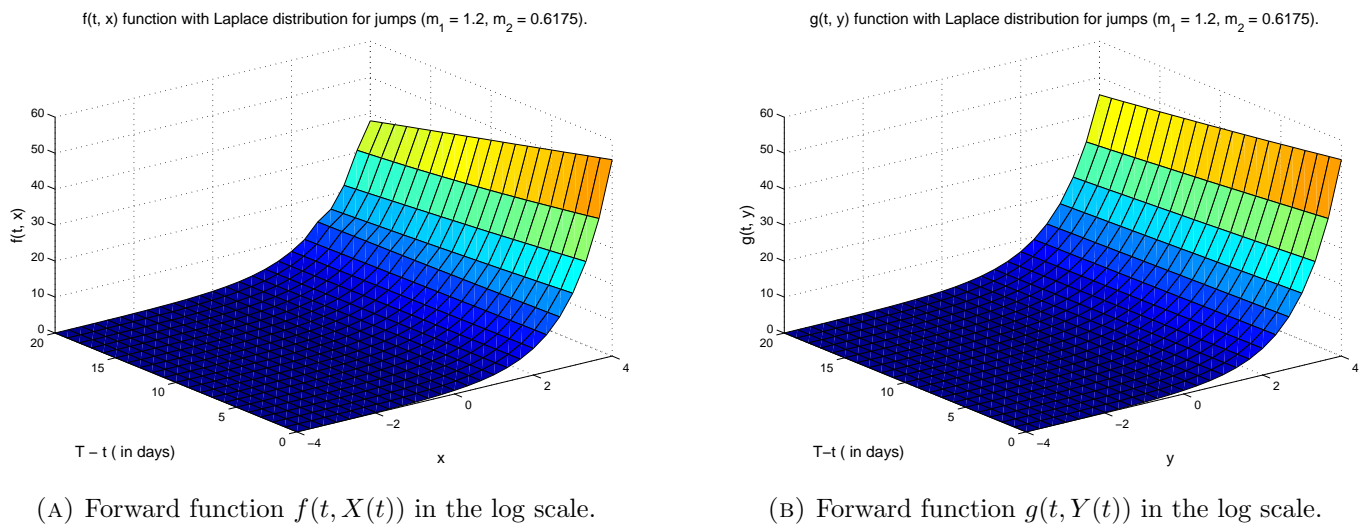


FIGURE 5. Laplace distribution assumed for jumps with location $m_1 = 1.2$ and scale $m_2 = 0.6175$. Time to maturity $T - t = 20$ days.

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